

The Entanglement Hamiltonian, SSD, and others in one-dimensional critical systems

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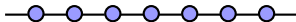
[In collaboration with
Xueda Wen (UIUC -> MIT)
and Andreas Ludwig (UCSB)]

[Closely related works: Ishibashi-Tada (15-16), Okunishi (16)]

Hamiltonians in CFT

- Let's start from a Hamiltonian of (1+1)d CFTs;

On a lattice (chain), it would look like:
$$H = \sum_i h_{i,i+1}$$



- Deformed evolution operator:
$$H[f] = \sum_i f\left(\frac{x_i + x_{i+1}}{2}\right) h_{i,i+1}$$

envelope function

- E.g. Entanglement Hamiltonian:
$$f(x) = \frac{R^2 - x^2}{2R}$$

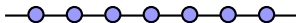
- E.g. Sine-square deformation (SSD):
$$f(x) = \sin^2 \frac{\pi x}{L}$$

[Gendiar-Krcmar-Nishino (2009) ...]

- Other applications: Rainbow chain, inhomogenous systems, quantum energy inequalities, etc.

Entanglement Hamiltonians in CFT

- Reduced density matrix for a finite interval $[-R, R]$ by tracing out its complement



$$\rho_A = \text{Tr}_{\bar{A}} |GS\rangle\langle GS|$$

- The entanglement Hamiltonian: $\rho_A \propto \exp(-H_E)$
- H_E is generically complicated.

- For CFT,

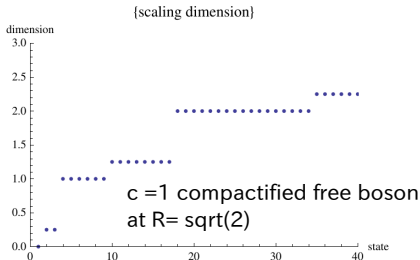
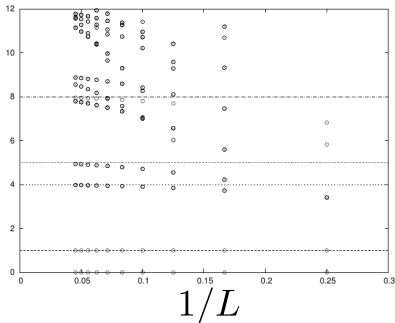
$$H_E = \int dx \frac{(R^2 - x^2)}{2R} \mathcal{H}(x)$$

- Finite size scaling: the spectrum of (boundary) CFT: Equidistance energy levels, and $1/\log(R)$ scaling

- For a given tower of states, all levels are equally spaced (with degeneracy, which depends on details of the theory)

$$H = \sum_i (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y)$$

$$(E - E_{GS}) / (E_1 - E_{GS})$$



- Level spacing scales as $1/L$

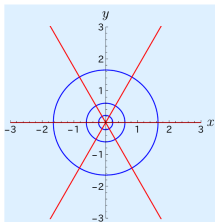
Strategy

- We will focus on critical systems (points) described by CFT.
- We will discuss types of "deformations" $H[f] = \int dx f(x)\mathcal{H}$ generated by various conformal maps.
- Put differently, we are interested in deformations which we can "undo" by conformal maps.
- Will discuss spectral properties (finite size scaling) of $H[f]$

Warm-up

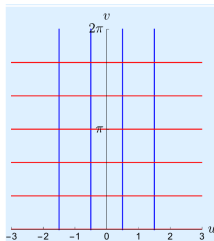
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$$z = x + iy$$



$$w(z) = \log z$$

$$w = u + iv$$



- CFT on the plane \leftrightarrow CFT on a cylinder (quantum lattice model on a circle)

- Radial evolution (Dilatation) \leftrightarrow Hamiltonian (1/L scaling)

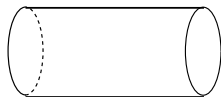
- Angular evolution ("Rindler" or "Modular" or "Entanglement" Hamiltonian) \leftrightarrow Hamiltonian with boundary

Finite size scaling of CFT [Cardy]

- CFT on a cylinder of circumference L :

$$\tilde{H} = \frac{1}{2\pi} \int_0^L dv \tilde{T}_{uu}(u_0, v)$$

$$\tilde{H} = \frac{1}{2\pi} \oint_{C_w} dw \tilde{T}(w) + (\text{anti-hol})$$



$$\tilde{T}_{uu}(w) = \tilde{T}(w) + \tilde{\bar{T}}(\bar{w})$$

- Conformal map: cylinder \rightarrow plane

$$w = \frac{L}{2\pi} \log z.$$

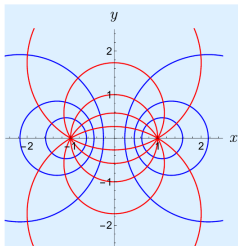
$$\tilde{T}(w) = \left(\frac{2\pi}{L}\right)^2 \left[z^2 T(z) - \frac{c}{24} \right]$$

$$\tilde{H} = \frac{2\pi}{L} \left(L_0 + \bar{L}_0 - \frac{c}{24} \right)$$

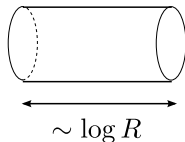
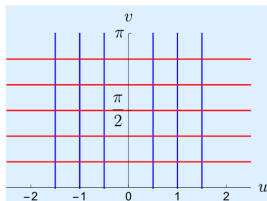
$$\oint_{C_w} dw \tilde{T}(w)$$

$$= \oint_{C_z} dz \frac{dw}{dz} \left(\frac{2\pi}{L}\right)^2 \left[z^2 T(z) - \frac{c}{24} \right]$$

$$= \oint_{C_z} dz \left(\frac{L}{2\pi}\right) \left[z T(z) - \frac{c}{24} \frac{1}{z} \right].$$



$$w(z) = \ln \frac{(z + R)}{(z - R)}$$



- Entanglement Hamiltonian on finite interval $[-R, R]$
 \leftrightarrow Hamiltonian with boundaries

- Transforming from strip to plane:

$$H = \int du T_{vv}(u, v_0 = \pi) = \int_{-R}^{+R} dx \frac{(x - R)(x + R)}{2R} T_{yy}(x, y = 0)$$

- Entanglement spec: $1/\text{Log}(R)$ scaling.

[See, e.g: Casini-Huerta-Myers (11), Cardy @ 2015 KITP conference]

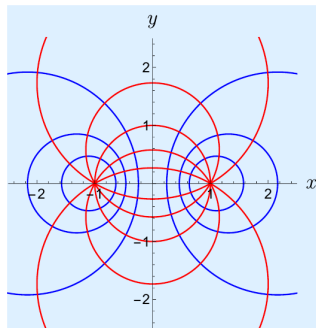
What is the evolution orthogonal to the evolution by Entanglement H ?

- Let's take a circle as a "space" (Cauchy surface)

$$\left(x + \frac{\cosh u_0}{\sinh u_0} R\right)^2 + y^2 = \frac{R^2}{(\sinh u_0)^2}$$

We have chosen: $r_0 := \frac{R}{\sinh u_0}$

Circumference: $L = 2\pi r_0$



- Evolution operator

$$H = \int_0^\pi dv T_{uu}(u_0, v) = r_0^2 \int_0^{2\pi} d\theta \frac{\cos \theta + \cosh u_0}{\sinh u_0} T_{rr}(r, \theta)$$

- "Regularized" version of the SSD:

$$H = \frac{L}{2\pi} \frac{1}{\sinh u_0} \int_0^L ds \left(\cos \frac{2\pi s}{L} + \cosh u_0 \right) T_{rr} \left(r = \frac{L}{2\pi}, \theta = \frac{2\pi s}{L} \right)$$

Summary

- The following Hamiltonian on a cylinder ($w = u_0 + iv$)

$$H = \frac{1}{2\pi} \int_0^L d\sigma \left[p - \frac{q}{2} \left(e^{-\frac{2\pi\xi}{L}} + e^{\frac{2\pi\xi}{L}} \right) \right] T_{\tau\tau}(\tau_0, \sigma)$$

$$p = r \cosh \lambda$$

$$q = r \sinh \lambda$$

can be mapped into the standard Hamiltonian on a cylinder by ($\xi = \tau_0 + i\sigma$):

$$\xi = \frac{L}{2\pi} \log \left[\frac{\cosh \frac{\pi w}{L} + \cosh \lambda}{\sinh \lambda} \right]$$

$$H = \frac{r}{2\pi} \int_0^L dv T_{uu}(u_0, v)$$

"Regularized" SSD

$$H = \frac{L}{2\pi} \frac{1}{\sinh u_0} \int_0^L ds \left(\cos \frac{2\pi s}{L} + \cosh u_0 \right) T_{rr} \left(r = \frac{L}{2\pi}, \theta = \frac{2\pi s}{L} \right)$$

- By construction, this operator has the spectrum of CFT on a circle with level spacing of order one.

- Define:

$$H_{\text{rSSD}} = \int_0^L ds \left(\cos \frac{2\pi s}{L} + \cosh u_0 \right) T_{rr} \left(\frac{L}{2\pi}, \frac{2\pi s}{L} \right)$$

- The envelope function:

$$\begin{aligned} f(s) &= \cos \left(\frac{2\pi s}{L} \right) + \cosh u_0 & f(s) &\xrightarrow{R \rightarrow 0} \cos \left(\frac{2\pi s}{L} \right) + 1 \\ &= \cos \left(\frac{2\pi s}{L} \right) + \sqrt{1 + \left(\frac{2\pi R}{L} \right)^2} & &= \cos^2 \left(\frac{\pi s}{L} \right) = \sin^2 \left[\frac{\pi}{L} \left(s - \frac{L}{2} \right) \right]. \end{aligned}$$

- "Regularized" version of the SSD:

R, the distance between vortices, is the regularization parameter.

- Scaling:

(i) fix u_0 , change R --> 1/L scaling

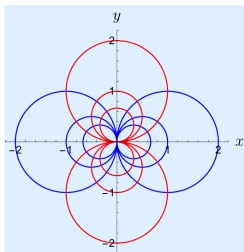
$$\sim \frac{\sinh u_0}{L} = \frac{1}{2\pi} \frac{(\sinh u_0)^2}{R} \sim \frac{1}{R}.$$

(ii) fix R, change u_0 --> 1/L² scaling

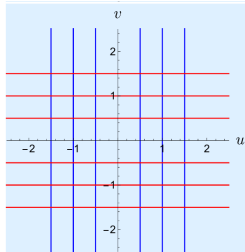
$$\sim \frac{\sinh u_0}{L} = \frac{1}{2\pi} \frac{(\sinh u_0)^2}{R} \sim \frac{1}{L^2}.$$

The dipolar limit

- Can take the dipolar limit $R \rightarrow 0$ rSSD \rightarrow SSD:



$$w(z) = 1/z$$



- In the dipolar limit, the w-plane (u-v plane) is an infinite plane

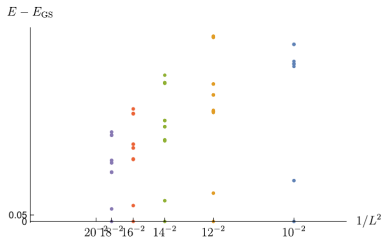
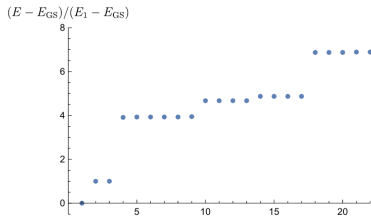
--> Infinite system length limit, continuum spectrum [Ishibashi-Tada (15,16)]

$$\begin{aligned} H &= \int_{-\infty}^{+\infty} dv T_{uu}(u_0, v) = 4r_0^3 \int_0^{2\pi} d\phi \sin^2(\phi/2) T_{rr}(r_0, \theta) \\ &= \frac{L^2}{\pi^2} \int_0^L ds \sin^2\left(\frac{\pi s}{L}\right) T_{rr}\left(\frac{L}{2\pi}, \frac{2\pi s}{L}\right) \end{aligned}$$

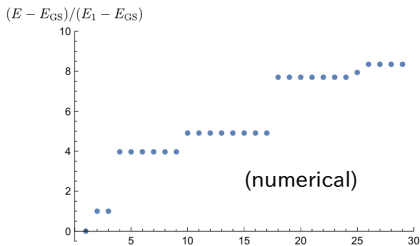
- The prefactor L^2 is indicative of the $1/L^2$ scaling seen in numerics.

Numerics (rSSD)

$$H = \sum_i (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y)$$



Regularized SSD spectrum

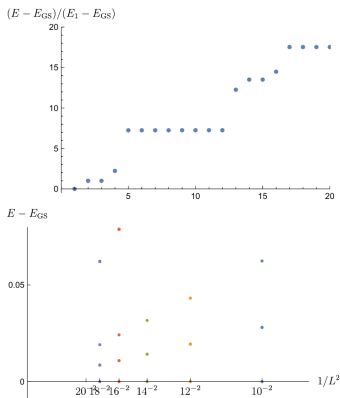


Physical spectrum (PBC)

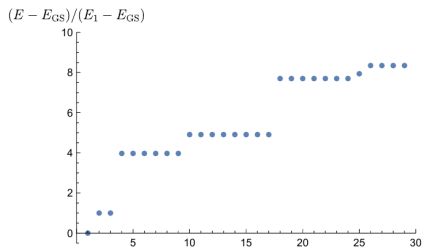
$$H = \sum_i f(x_i + 1/2) h_{i,i+1}$$

$$f(x) = \cos \frac{2\pi x}{L} + \sqrt{1 + \left(\frac{2\pi R}{L}\right)^2}$$

Numerics (SSD)



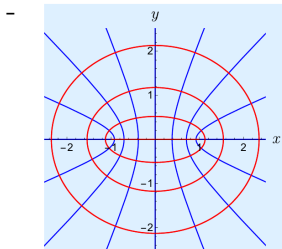
SSD



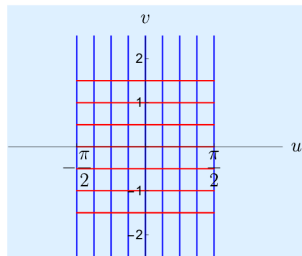
Physical spectrum (PBC)

SSD spectrum does not much physical spectrum,
 $1/L^2$ scaling

Other examples



$$z = \sin(w)$$



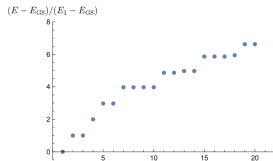
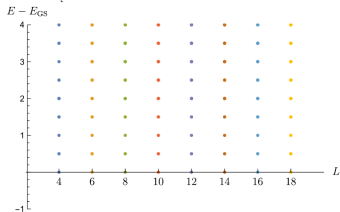
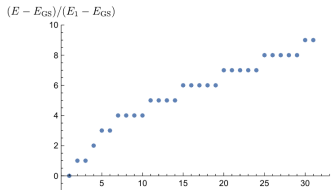
infinite stripe

- Engineering conformal map and evolution operator

$$-\pi/2 < u < \pi/2$$
$$-\infty < v < +\infty$$

- "square root deformation" $\tilde{H} = \int_{-R+\epsilon}^{+R-\epsilon} dx \sqrt{R^2 - x^2} T_{yy}$

- Known in the context of "perfect state transfer"
(Thanks: Hosho Katsura)



Physical spectrum (OBC)

"Square root" deformation

Summary

- Setup a general discussion of "deformed" Hamiltonians in CFTs
- Proposed a "regularized" version of SSD (rSSD).
- Original SSD can be viewed as a "singular" limit of rSSD
- Spectrum of rSSD is easy to understand.
Shed light on $1/L^2$ scaling of SSD.